

# General Relativity: Example Sheet 2

David Tong, October 2019

1. Consider a change of basis  $\tilde{e}_\mu = (A^{-1})^\nu{}_\mu e_\nu$ . Show that the components of a connection in the new basis are related to its components in the old basis by

$$\tilde{\Gamma}^\rho{}_{\mu\nu} = A^\rho{}_\lambda (A^{-1})^\sigma{}_\mu [(A^{-1})^\tau{}_\nu \Gamma^\lambda{}_{\sigma\tau} + e_\sigma((A^{-1})^\lambda{}_\nu)]$$

Show further that the difference of two connections,  $(\Gamma_1)^\rho{}_{\mu\nu} - (\Gamma_2)^\rho{}_{\mu\nu}$ , transforms as a tensor.

2. Let  $\nabla$  be a connection that is not torsion-free. Let  $T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$  where  $X$  and  $Y$  are vector fields. Show that this defines a  $(1, 2)$  tensor field  $T$ . This is called the *torsion tensor*. Show that, for any function  $f$ ,

$$2\nabla_{[\mu} \nabla_{\nu]} f = -T^\rho{}_{\mu\nu} \nabla_\rho f$$

3. Let  $\nabla$  be a torsion-free connection. Derive the analogue of the Ricci identity for a 1-form  $\omega$ ,

$$2\nabla_{[\mu} \nabla_{\nu]} \omega_\rho = -R^\sigma{}_{\rho\mu\nu} \omega_\sigma$$

4. The Riemann tensor constructed from the Levi-Civita connection obeys the Bianchi identity  $R^\mu{}_{\nu[\rho\sigma;\lambda]} = 0$ . Use this fact to derive the contracted Bianchi identity  $G^\mu{}_{\nu;\mu} = 0$  where  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu}$  is the Einstein tensor.

5. A vector field  $Y$  is parallelly propagated (with respect to the Levi-Civita connection) along an affinely parameterized geodesic with tangent vector  $X$  in a Riemannian manifold. Show that the magnitudes of the vectors  $X, Y$  and the angle between them are constant along the geodesic.

On the unit sphere a unit vector  $Y$  is initially tangent to the line  $\phi = 0$  at a point on the equator. It is then moved by parallel propagation first along the equator to the point  $\phi = \phi_0$ , from there along the line  $\phi = \phi_0$  to the North pole, and then back along the line  $\phi = 0$  to its original position. By how much has it changed, and why?

**6\***. The *Reissner-Nordstrom* solution of the Einstein-Maxwell equations has metric

$$ds^2 = -f(r)^2 dt^2 + f(r)^{-2} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

with

$$f(r)^2 = 1 - \frac{2M}{r} + \frac{Q^2 + P^2}{r^2}$$

and a Maxwell field strength  $F = dA$ , with

$$A = -\frac{Q}{r} dt - P \cos \theta d\phi$$

where  $M, P, Q$  are constants.  $M$  can be interpreted as the total mass of this spacetime. Assume that  $(t, r, \theta, \phi)$  is a right handed coordinate chart. Show that

$$\frac{1}{4\pi} \int_{\mathbf{S}_\infty^2} \star F = Q \quad \text{and} \quad \frac{1}{4\pi} \int_{\mathbf{S}_\infty^2} F = P \quad (1)$$

where  $\mathbf{S}_\infty^2$  is a sphere at  $r = \infty$  on a surface of constant  $t$ . What is the physical interpretation of  $Q$  and  $P$ ?

**7.** In Q5 of examples sheet 1, we showed that

$$\begin{aligned} (\mathcal{L}_X \omega)_\mu &= X^\nu \partial_\nu \omega_\mu + \omega_\nu \partial_\mu X^\nu \\ (\mathcal{L}_X g)_{\mu\nu} &= X^\rho \partial_\rho g_{\mu\nu} + g_{\mu\rho} \partial_\nu X^\rho + g_{\rho\nu} \partial_\mu X^\rho \end{aligned}$$

Use normal coordinates to argue that one can replace partial derivatives with covariant derivatives to obtain the basis-independent results

$$\begin{aligned} (\mathcal{L}_X \omega)_\mu &= X^\nu \nabla_\nu \omega_\mu + \omega_\nu \nabla_\mu X^\nu \\ (\mathcal{L}_X g)_{\mu\nu} &= \nabla_\mu X_\nu + \nabla_\nu X_\mu \end{aligned}$$

with  $\nabla$  is the Levi-Civita connection.

**8.** How many independent components does the Riemann tensor (of the Levi-Civita connection) have in two, three and four dimensions? Show that in two dimensions

$$R_{\mu\nu\rho\sigma} = \frac{1}{2} R (g_{\mu\rho} g_{\nu\sigma} - g_{\mu\sigma} g_{\nu\rho}).$$

Discuss the implications for general relativity in two spacetime dimensions.

9. In a  $d$ -dimensional spacetime, define a tensor

$$C_{\mu\nu\rho\sigma} = R_{\mu\nu\rho\sigma} + \alpha(R_{\mu\rho}g_{\nu\sigma} + R_{\nu\sigma}g_{\mu\rho} - R_{\mu\sigma}g_{\nu\rho} - R_{\nu\rho}g_{\mu\sigma}) + \beta R(g_{\mu\rho}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\rho})$$

where  $\alpha$  and  $\beta$  are constants. Show that  $C_{\mu\nu\rho\sigma}$  has the same symmetries as  $R_{\mu\nu\rho\sigma}$ .

What values of  $\alpha$  and  $\beta$  give  $C^\mu{}_{\nu\mu\sigma} = 0$ ? Determine them. With this extra condition  $C_{\mu\nu\rho\sigma}$  is called the *Weyl tensor*. Show that it vanishes if  $d = 2, 3$ .

Setting  $d = 4$ , how many independent components do  $R_{\mu\nu}$  and  $C_{\mu\nu\rho\sigma}$  have? Show that in vacuum

$$\nabla^\mu C_{\mu\nu\rho\sigma} = 0.$$

What does the Weyl tensor represent physically?

10. [Optional] Use the Bianchi identity to derive the *Penrose equation* for a vacuum spacetime

$$\nabla^\lambda \nabla_\lambda R_{\mu\nu\rho\sigma} = 2R^\kappa{}_{\mu\lambda\sigma} R^\lambda{}_{\rho\kappa\nu} - 2R^\kappa{}_{\nu\lambda\sigma} R^\lambda{}_{\rho\kappa\mu} - R^\kappa{}_{\lambda\sigma\rho} R^\lambda{}_{\kappa\mu\nu}$$

11\*. Consider metrics of the form

$$ds^2 = -f(r)^2 dt^2 + f(r)^{-2} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

Use the action for a test particle to write down the geodesic equations in this metric, and hence extract the Christoffel symbols in coordinates  $(t, r, \theta, \phi)$ .

Use a basis of vierbeins to determine the curvature 2-form, and hence the components of the Riemann tensor in coordinates  $(t, r, \theta, \phi)$ .